

# A Semi-Theoretical Study on the Effect of Rough Beds on the Conjugate Depth Ratio in Hydraulic Jumps

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## ABSTRACT

This study examines the effects of bed roughness on hydraulic jump characteristics in rectangular composite channels. Experiments were conducted in a laboratory flume with two roughness zones: the main channel and the floodplain, using uniform plastic granules to simulate different roughness conditions. The results show that increased bed roughness significantly reduces the subsequent depth ratio. A semi-theoretical dimensionless approach was developed to analyze the hydraulic jump, incorporating key parameters, such as the incident Froude number, the resistance coefficient, the expansion ratio, the shape factor, and the conjugate depth ratio. This approach, based on Euler's theory applied between two sections of the hydraulic jump, allows the conjugate depth ratio to be predicted using the flow rate and the tested absolute roughness values. The experimental validation confirms the accuracy of this model in estimating the conjugate depth ratio under varying roughness conditions. The developed approach provides a practical tool for engineers to optimize hydraulic structures in channels with variable roughness and zero slope, improving flow regulation and reducing erosive forces.

*Keywords-hydraulic jump; bed roughness; conjugate depth ratio; rectangular compound channel; rough bed*

## I. INTRODUCTION

The hydraulic jump, a rapid transition from supercritical to subcritical flow, plays a critical role in energy dissipation and scour prevention in hydraulic structures [1, 2]. Classical theories have established fundamental relationships for jump characteristics in smooth rectangular channels [3], but real-world applications involving complex geometries and rough bed conditions require advanced analysis of turbulence dynamics and energy dissipation [4, 5]. Authors in [6, 7] show that rough beds enhance energy dissipation and reduce the subsequent depth ratios by altering turbulence structures. For

example, corrugated or pebbled beds induce secondary flows, increase shear stresses, and stabilize jump positions [8, 9]. Submerged jumps over rough transition beds further reveal universal scaling laws for turbulence intensities [10, 11], highlighting the interaction between the roughness elements and flow regimes [12, 13]. Compound channels with distinct main and floodplain beds introduce additional complexity due to lateral variations in geometry and roughness. While authors in [14] examined hydraulic jumps in straight compound channels, authors in [15, 16] showed that roughened stilling basins reduce jump lengths by up to 40%. Research on triangular corrugated beds [17] emphasizes the role of bed

orientation in flow separation, while compound roughness configurations in inclined channels alter the roll length and energy loss [18, 19]. However, these studies often overlook roughness interactions in multi-bed systems under zero-slope conditions, despite their prevalence in urban drainage networks [20]. Authors in [21, 22] examined the effects of roughness on bounce characteristics, confirming that bed texture significantly affects conjugate depth ratios and turbulence patterns in rectangular channels under varying flow conditions. The objective of this study was to develop a semi-theoretical mathematical model, a hybrid approach that integrates fundamental theoretical equations (e.g., Euler's equations or momentum conservation principles) with experimentally derived relationships, to analyze hydraulic jump dynamics in rectangular compound channels. The research began by applying Euler's equations between two cross-sections of a stabilized hydraulic jump, setting up dimensionless relationships between key parameters, such as incident Froude number ( $F_r$ ), conjugate depth ratio ( $Y$ ), resistance coefficient ( $C_r$ ), channel expansion ratio ( $\beta$ ), and shape factor ( $\tau$ ). Although the theoretical framework provided the foundation, critical unknowns such as  $C_r$ , which depends on channel roughness and flow conditions, were resolved empirically through laboratory experiments. A specialized test rig allowed comprehensive measurements under controlled conditions, with four uniformly distributed plastic roughness configurations ( $\varepsilon = 6, 8, 10, 12$  mm) ensuring homogeneity, unlike previous studies using irregular surfaces [23, 24]. By explicitly correlating  $C_r$  with measurable roughness ( $\varepsilon$ ), the model combined universal physical laws with site-specific empirical data, yielding highly accurate dimensionless equations for predicting  $Y$  using variables, such as flow rate and  $\varepsilon$ . This semi-theoretical approach differs from fully theoretical models (which rely solely on closed-form equations) and purely empirical models (which lack physical principles), and offers a robust tool for predicting hydraulic jump behavior in different flow scenarios, thereby advancing methods in mixed-channel hydraulics.

## II. EXPERIMENTAL PROTOCOL

The study was conducted using a closed-loop hydraulic flume system at the Laboratory for Exploitation and Valorization of Natural Resources in Arid Zones, University of Ouargla, as shown in Figure 1. The primary test facility consisted of a 10-meter-long, horizontal (zero slope), rectangular composite channel with transparent plexiglas side walls for flow visualization. The channel was divided into a 4-meter experimental section with a compound cross-section consisting of a minor bed (14.4 cm wide) and a major bed (25 cm wide), both 15.5 cm high. A 150 mm diameter circular pipe supplied water from an axial pump (maximum capacity: 55.55 L/s) to an upstream closed metal box containing a flat orifice plate to induce a controlled torrential flow. The orifice's adjustable opening corresponded to the initial hydraulic head ( $h_1$ ), which was controlled by a downstream sluice. A rectangular sharp-crested weir, installed at the channel's downstream end, enabled direct flow measurement. The experiments were conducted using four initial heights ( $h_1 = 2.5$  cm, 3 cm, 3.5 cm, and 4 cm), selected to achieve practical Froude number ranges ( $F_r \approx 2 - 9$ ) relevant to real-world

scenarios, like spillways and stilling basins, ensuring stable jump formation. The experiments were performed with four roughness values ( $\varepsilon = 06$  mm, 08 mm, 10 mm, and 12 mm), chosen to reflect industrial and natural analogs (e.g., concrete channels and rock-bedded streams) were utilized, with  $\varepsilon/B$  ratios maintaining flow homogeneity.

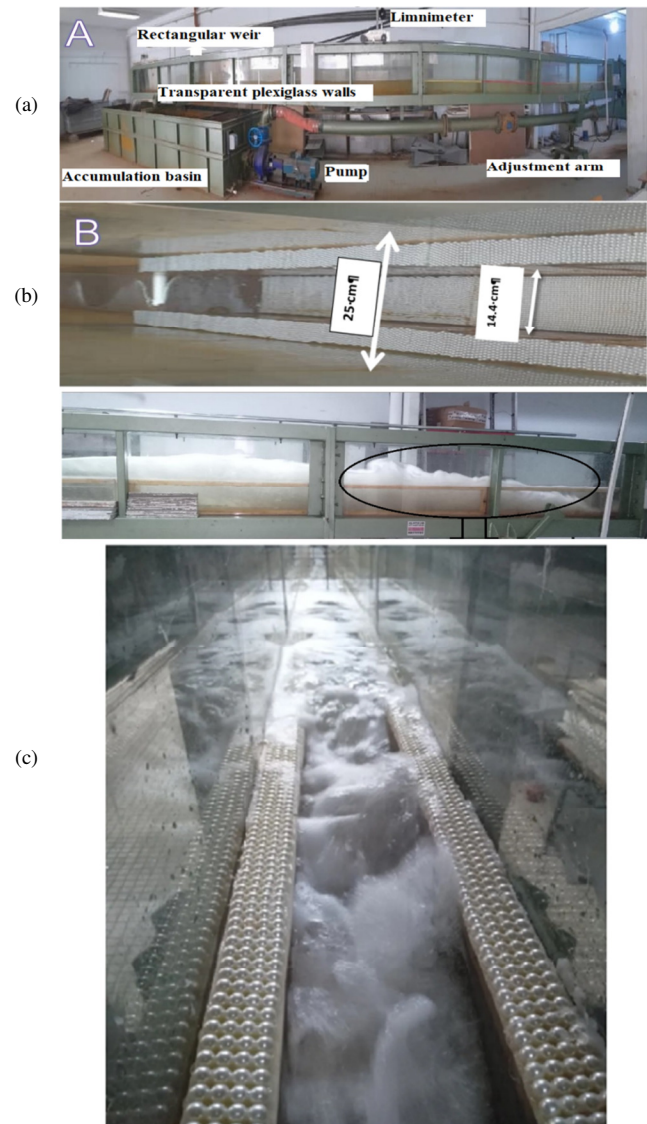


Fig. 1. (a) The experimental channel used for studying the hydraulic jump, (b) a cross-section of the compound channel showing the dimensions of the main bed and the secondary bed, (c) hydraulic jump formation in a rectangular compound channel with small and large rough beds.

Plastic pellets with a central hole were employed as roughness elements, securely attached with a rigid fishing line and arranged linearly along the channel bed. These elements were glued with a robust adhesive to prevent erosion and detachment during torrential flow, offering durability for multiple uses. Materials, like paper, glass, sheet metal, or gravel, were selected over for their rigidity and resistance to high-flow conditions, as presented in Figure 2. In addition, the

study examined 26 adjustable thresholds (2.5 cm –21 cm) composed of thin metal plates (1 mm –2 mm thick) anchored to the channel bed in 1 cm increments. These thresholds were utilized to vary tailwater levels and cover subcritical to supercritical flow transitions. The influence of these thresholds on jump control was assessed. A schematic representation of the model is provided in Figure 3, including the channel layout, mat positions, and measurement points.

III. THEORETICAL ANALYSIS

The objective of this section is to establish a theoretical dependence governing the transition of the flow from the torrential (supercritical) regime to the fluvial (subcritical) regime in a rough compound channel, using Euler's theory. The desired dimensionless function is expressed as  $f(Y, F_r, C_r, \beta, \tau) = 0$ , where  $Y$  is the ratio of flow depths ( $h_2/h_1$ ),  $F_r$  is the upstream Froude number,  $\tau = h_1/h_c$  is the shape ratio,  $\beta = b/B$  is the channel width expansion ratio,  $C_r$  is the resistance parameter accounting for roughness effects. The momentum balance between the initial (upstream) and final (downstream) sections of the hydraulic jump is:

$$\rho Q(v_1 - v_2) = F_1 - F_2 - F_r \tag{1}$$

where  $F_1$  and  $F_2$  are the hydrostatic forces at the upstream and downstream sections, and  $F_R$  is the friction force due to roughness, as depicted in Figure 4.

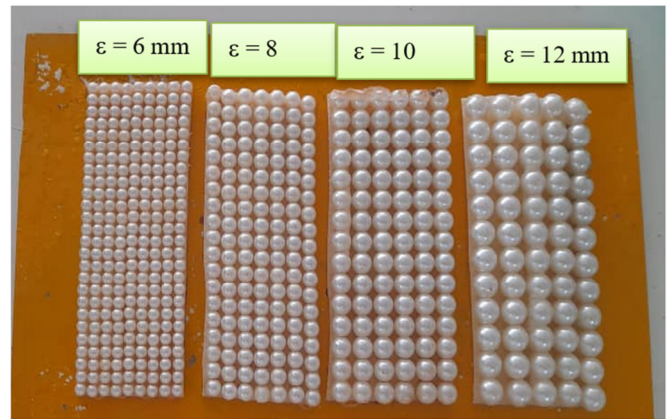


Fig. 2. Mats with varying roughness used in hydraulic jump experiments.

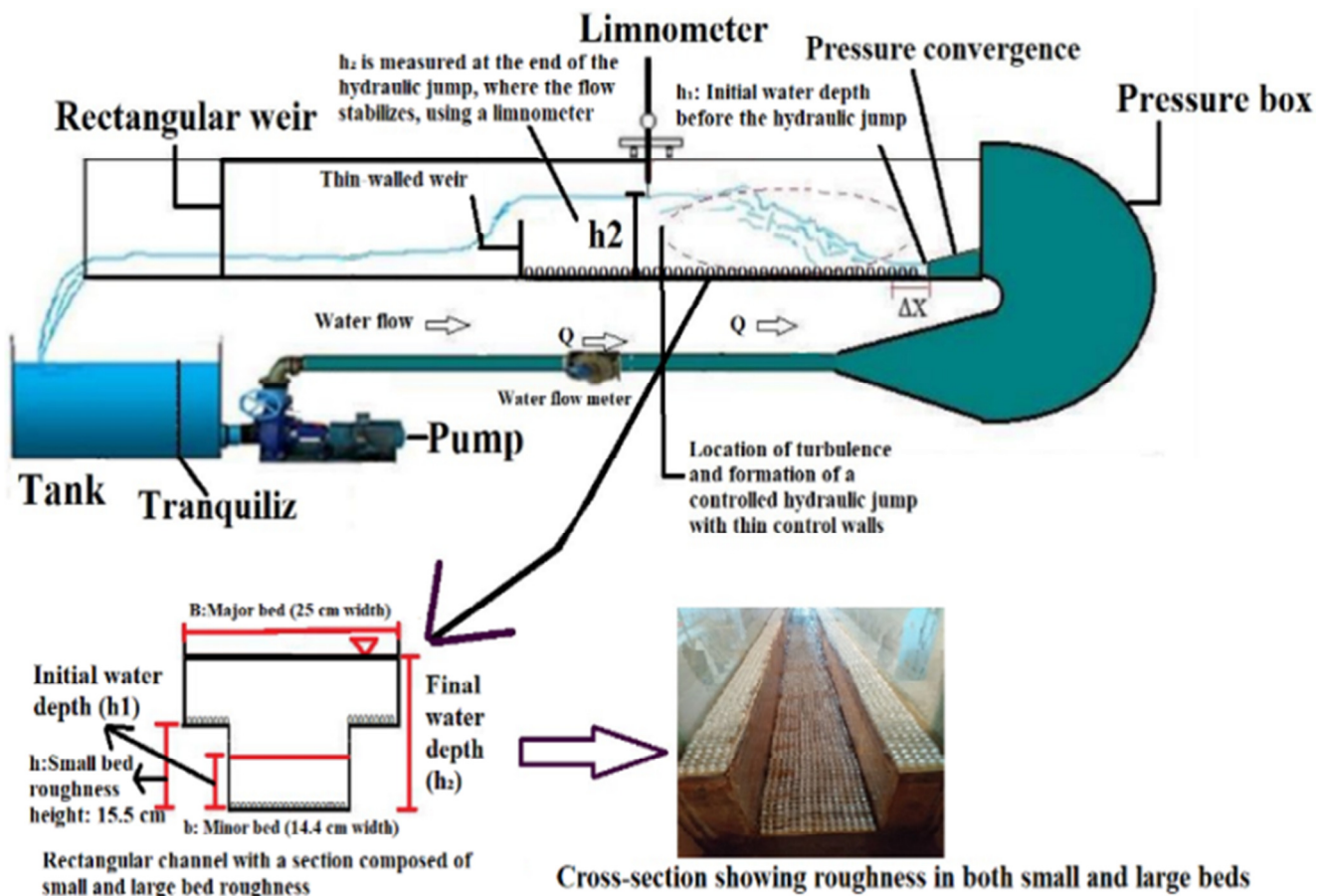


Fig. 3. Longitudinal section of the experimental model showing the channel layout, mat positions, and measurement points.

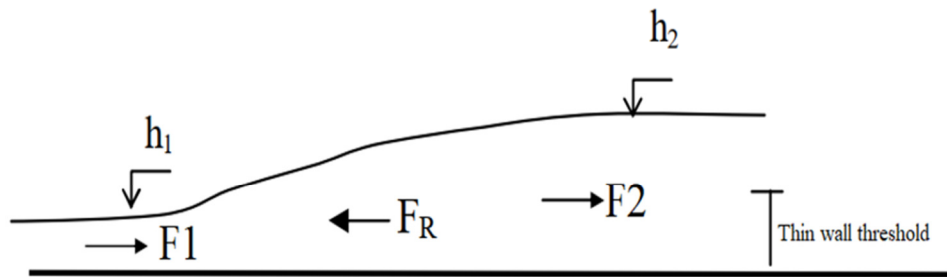


Fig. 4. Schematic representation of a hydraulic jump over a thin-wall threshold.

The initial and final sections are:

$$A_1 = bh_1 \tag{2}$$

$$A_2 = B(h_2 - h) + bh \tag{3}$$

The magnitude of hydrostatic forces is related to the water depth and channel width:

$$F_1 = \rho Q \left( \frac{h_1}{2} \right) bh_1 \tag{4}$$

$$F_2 = \rho Q(h_2 - h/2)bh + \rho Q[(h_2 - h)/2]B(h_2 - h) \tag{5}$$

where,  $h_1$  and  $h_2$  are the depths,  $b$  and  $B$  are the widths,  $\rho$  is water density, and  $g$  is the gravitational acceleration. Roughness resistance is:

$$F_r = C_r \rho g L_j \left( \frac{V^2}{2g} \right) p \tag{6}$$

where  $C$  is the roughness coefficient,  $L_j$  is the jump length,  $p=B$  is the wetted perimeter, and  $V$  is the flow velocity. The adjusted equation becomes:

$$\rho Q^2 \left( \frac{1}{bh_1} - \frac{1}{B(h_2-h)+bh} \right) = -\bar{\omega} \left( \frac{h_1}{2} \right) bh_1 + \bar{\omega} \left( h_2 - \frac{h}{2} \right) bh + \bar{\omega} \left[ \frac{(h_2-h)}{2} \right] B(h_2 - h) \tag{7}$$

The final modified equation becomes:

$$F_r^2 \left( 1 - \frac{C_r}{\beta} \right) = \frac{-1 + \frac{(2Y-\frac{1}{\tau}) + (\frac{Y-\frac{1}{\tau}}{\beta})^2}{2}}{Y - \frac{(1-\beta)}{\tau}} + \frac{\beta F_r^2}{Y - \frac{(1-\beta)}{\tau}} \tag{8}$$

where the incident Froude number is given by  $F_r = \frac{Q}{\sqrt{b^2 h_1^3 g}}$ ,

while the channel expansion ratio is expressed as  $\beta = \frac{b}{B}$ , the conjugate depth ratio is  $Y = \frac{h_2}{h_1}$ , whereas the shape factor is defined as  $\tau = \frac{h_1}{h}$ , and the friction coefficient is  $C_r = \frac{L_j}{h_1}$ .

#### IV. EXPERIMENTAL STUDY OF THE SEMI-THEORETICAL RELATION

This section evaluates the friction coefficient  $C_r$  through laboratory experiments. The study's objective is to ascertain the variation of  $C_r$  with the Froude number ( $F_r$ ) and surface roughness  $\varepsilon$ . The semi-theoretical relation is subjected to experimental testing and expressed by (8). To facilitate the experimental analysis, the expression  $f(Y)$  is defined as:

$$f(Y) = F_r^2 \left( 1 - \frac{C_r}{\beta} \right) \tag{9}$$

Substituting this into (8):

$$f(Y) = \frac{-1 + \frac{(2Y-\frac{1}{\tau}) + (\frac{Y-\frac{1}{\tau}}{\beta})^2}{2}}{Y - \frac{(1-\beta)}{\tau}} + \frac{\beta F_r^2}{Y - \frac{(1-\beta)}{\tau}} \tag{10}$$

The variation of  $f(Y)$  with  $F_r^2$  is analyzed for four roughness values:  $\varepsilon = 06$  mm,  $\varepsilon = 08$  mm,  $\varepsilon = 10$  mm, and  $\varepsilon = 12$  mm. These values are different levels of bed roughness that affect the coefficient of friction  $C_r$ , as shown in Figure 5.

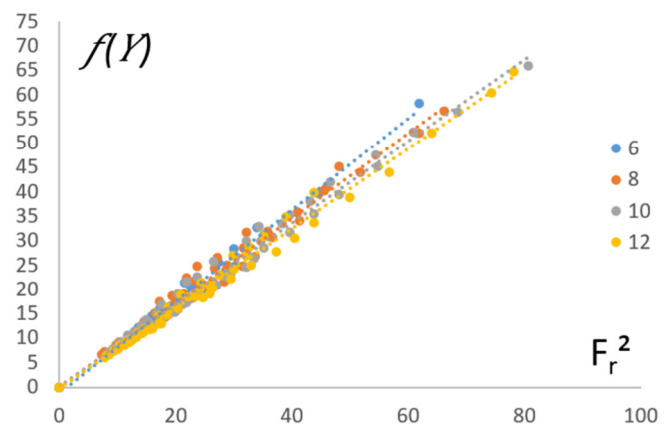


Fig. 5. Variation of  $f(Y) = F_r^2 \left( 1 - \frac{C_r}{\beta} \right)$  as a function of  $F_r^2$  with different bed roughness values.

A thorough examination of the experimental data collected from hydraulic jump measurements reveals that the distribution of data points for each pattern can be precisely represented by a linear relationship:

$$F_r^2 T = \frac{-1 + \frac{(2Y-\frac{1}{\tau}) + (\frac{Y-\frac{1}{\tau}}{\beta})^2}{2}}{Y - \frac{(1-\beta)}{\tau}} + \frac{\beta F_r^2}{Y - \frac{(1-\beta)}{\tau}} \tag{11}$$

where  $T$  is the factor  $\left( 1 - \frac{C_r}{\beta} \right)$ . A summary of the  $C_r$  coefficients is provided in Table I.

The data presented, suggest a positive correlation between the parameter  $C_r$  and the relative roughness ( $\varepsilon/B$ ). A statistical correlation analysis of the paired values ( $\varepsilon/B$ ,  $C_r$ ) gives a linear relationship that can be expressed by:  $C_r = 0.4166 (\varepsilon/B)$ , as

portrayed in Figure 6. By substituting the expression for  $C_r$  into (11), the semi-theoretical equation is transformed as:

$$F_r^2 \left( 1 - \frac{0.4166 (\frac{\epsilon}{B})}{\beta} \right) = \frac{-1 + \frac{(2Y - \frac{1}{\tau}) + (Y - \frac{1}{\tau})^2}{\tau}}{2} + \frac{\beta F_r^2}{Y - \frac{(1-\beta)}{\tau}} \quad (12)$$

The curve in Figure 7 presents a strong relationship between  $f(Y)$  and  $F_r^2 (1 - \frac{0.4166 (\epsilon/B)}{\beta})$ . The data points are distributed in a close pattern, suggesting a substantial degree of agreement between the experimental data and the theoretical predictions. This distribution confirms the accuracy and reliability of the measurements and shows that the actual values of  $f(Y)$  agree well with the predicted values calculated using the given mathematical relationship.

TABLE I. THE VALUES OF THE CR

Roughness (mm)	Relative roughness ( $\epsilon/B$ )	Parameter ( $C_r$ )	$R^2$
6	0.024	0.0688	0.9835
8	0.032	0.0792	0.9804
10	0.04	0.0924	0.9886
12	0.048	0.1095	0.9892
0 (smooth)	0	0	1

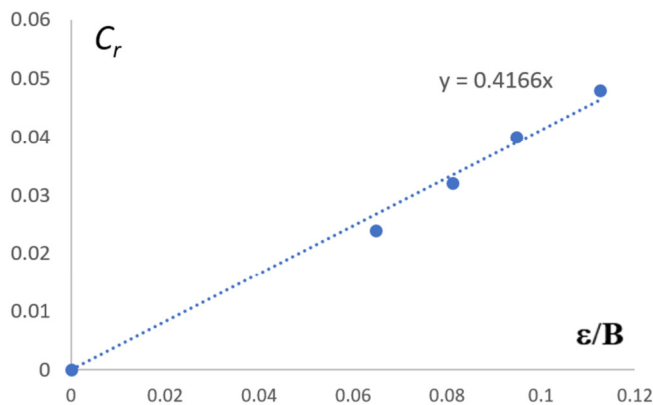


Fig. 6. Variation of  $C_r$  as a function of  $\epsilon/B$ .

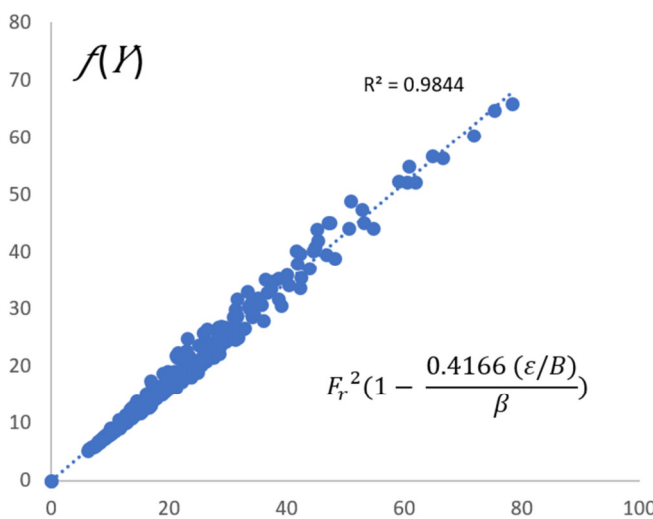


Fig. 7. Variation of  $f(Y)$  with respect to  $F_r^2 (1 - \frac{0.4166 (\epsilon/B)}{\beta})$ .

The application of the Buckingham  $\pi$  theorem resulted in a general form  $Y = a F_r^m (1 - \frac{0.4166 (\epsilon/B)}{\beta})^n$ , displayed in Figure 8. The unknown coefficients ( $a, m, n$ ) were determined through nonlinear regression using experimental measurements. The least squares method was then employed to optimize these parameters:

$$Y = 1.329 F_r^{0.986} (1 - \frac{0.4166 (\epsilon/B)}{\beta})^{1.430} \quad (13)$$

The graph shown in Figure 9 demonstrates the relationship between the experimental values ( $Y_{Experimental}$ ) and a mathematical model based on an empirical equation. The data points exhibit a high degree of alignment with the trend line, suggesting a strong correlation between the predicted and experimental values.

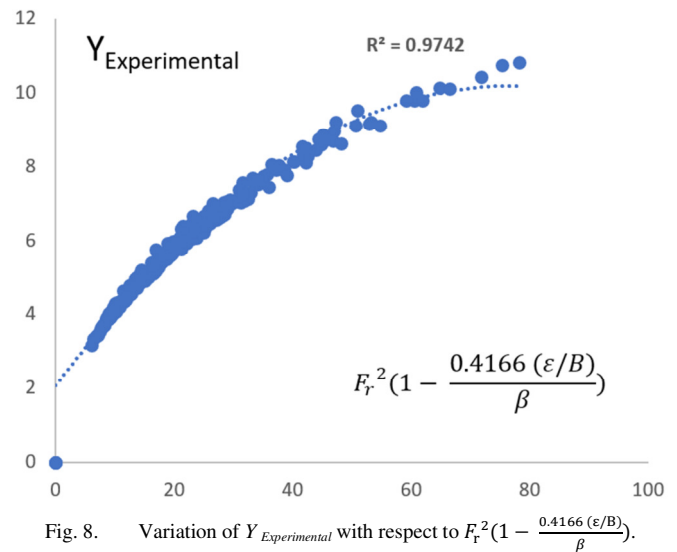


Fig. 8. Variation of  $Y_{Experimental}$  with respect to  $F_r^2 (1 - \frac{0.4166 (\epsilon/B)}{\beta})$ .

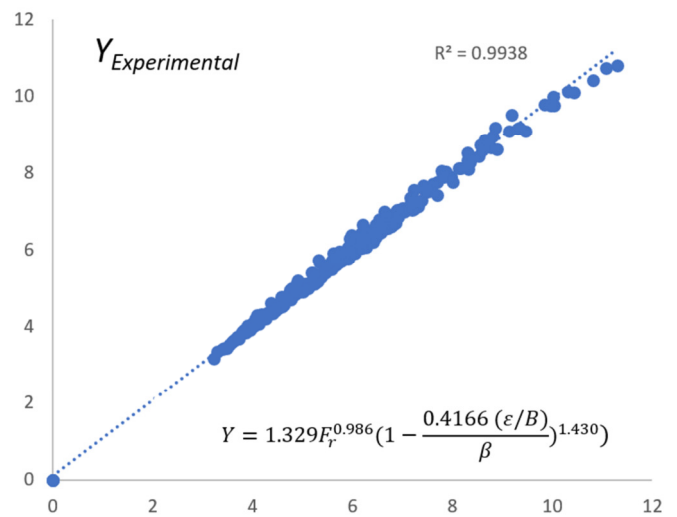


Fig. 9. Variation of  $Y_{Experimental}$  with respect to  $Y = 1.329 F_r^{0.986} (1 - \frac{0.4166 (\epsilon/B)}{\beta})^{1.430}$ .

## V. CONCLUSIONS

This study shows that bed roughness significantly influences hydraulic jump characteristics in rectangular compound channels. The experimental results reveal that increasing roughness reduces the sequent depth ratio, following the empirical trend  $Y = 1.329F_r^{0.986} \left(1 - \frac{0.4166(\epsilon/B)}{\beta}\right)^{1.430}$ , indicating that the sequent depth ratio ( $Y$ ) decreases as the friction coefficient  $C_r$  increases. A semi-theoretical, dimensionless model was developed based on Euler's equations applied between two sections of the hydraulic jump. The model incorporates key hydraulic parameters, including the incident Froude number, resistance coefficient, expansion ratio, and absolute roughness. This enables an accurate prediction of the conjugate depth ratio. The experimental validation confirms that the model reliably captures the effects of roughness variations on hydraulic jump behavior. The findings offer a pragmatic approach for optimizing hydraulic structures, particularly in energy dissipation systems, flood control channels, and stilling basins, underscoring the significance of roughness design in stabilizing hydraulic jumps and mitigating structural erosion.

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